



## (Magnetic) Small Angle Neutron Scattering Theory, Instrumentation and Applications



MLZ is a cooperation between:



Helmholtz-Zentrum Geesthacht Zentrum für Material- und Küstenforschung





Acknowledgments



### Research at large scale facilities: A team effort!

| P. Böni<br>C. Pfleiderer<br>F. Jonietz<br>T. Adams<br>J. Kindervater<br>T. Reimann<br>A. Backs<br>D. Mettus<br>F. Haslbeck | SCES TUM<br>E21/E51 | A. Heinemann<br>S. Busch<br>A. Beldowski<br>O. Listing<br>D. Heims<br>J. Borchers<br>G. Musilak<br>D. Siemers<br>J. Hedde<br>J. Hübscher | HZG        |
|--|---------------------|--|------------|
| A. Rosch   | Köln                | S. Demirdis  | JCNS       |
| M. Garst<br>B. Wolba   | Karlsruhe           | V. Pipich  |            |
| A. Wilhelm<br>S. Semecky<br>M. Schulz<br>K. Zeitelhack<br>I. Defendi<br>W. Petry<br>U. Reinecke<br>R. Gilles<br>X. Brems   | FRM II              | E. M. Forgan<br>A. Holmes<br>E. Blackburn<br>M. Laver  | Birmingham |
|  |                     | H. Furukawa<br>M. Soda<br>N. Kagamida  | RIKEN      |
| M. Dembski-Villalta  |                     | U. Keiderling  | HZB        |
|  |                     | R. Cubitt  | ILL        |







- SANS Basic Concept & Theory
- SANS Instrumentation
- SANS Resolution & Intensity
- Applications & Examples of SANS:
- Soft Matter
- Hard Matter
- Magnetism of nanoscale materials
- Emergent nanoscale magnetic structures of strongly correlated electron systems (SCES)





SANS: Diffractometer specialized for small scattering angles

**FRM II** 

Heinz Maier-Leibnitz

#### Low **Q** small scattering angles $\sim 1 \text{ Å}^{-1}$ to $\sim 10^{-4} \text{ Å}^{-1}$ Large correlations in real space 20Å to 40.000Å









### A joint TUM/HZG project





# SANS vs. SAXS



### Properties of neutrons

Interaction with the nuclei (*strong* interaction, pointlike)

- Neutral particle, deep penetration (window materials for extreme environment)
- Isotope sensitivity
- No radiation damage to samples
- Sensitive to magnetism (spin1/2 particle)
- Energy and momentum match elementary excitations and interatomic distances of condensed matter
- Low brilliance (many orders of magntitude compared to X-rays)
- Brilliance cannot be scaled up easily



#### Table 1.2: Properties of the Neutron.

| Physical quantity        | Quantity               | Dimension                                 |
|--------------------------|------------------------|---|
| Mass                     | $1.675 \cdot 10^{-27}$ | kg  |
| Charge                   | 0                      | С   |
| Spin                     | 1/2                    | $\hbar$                                   |
| magn. dipol moment       | $\mu_n = -1.913\mu_K$  | $\mu_K = \frac{e\hbar}{2M_p c}$           |
| nuclear magneton         |                        | $1\mu_K = 0.505\cdot 10^{-23}~{ m erg/G}$ |
|                          |                        | $1\mu_K = 3.15\cdot 10^{-14}~{ m MeV/T}$  |
| life time (free neutron) | 886                    | S   |
| kinetic energy           | $E = \frac{1}{2}mv^2$  | meV                                       |





#### Forward scattering - Why is it useful?



 $\Rightarrow$  Lots of intensity scattered in forward direction



Scattering length density: "Average interaction potential" for neutrons

# Principle of Babinet

T Z

τ Ζ



Assume a general two phase system

**FRM II** 

Forschungs-Neutronenguelle

Heinz Maier-Leibnitz

$$\begin{aligned}
V &= V_1 + V_2 \\
\rho(r) &= \begin{cases}
\rho_1 & \text{in } V_1 \\
\rho_2 & \text{in } V_2
\end{aligned}$$

T T Z



SANS measures inhomogeneities of scattering length density

$$\frac{d\Sigma}{d\Omega}(\mathbf{q}) = \frac{1}{V}(\rho_1 - \rho_2)^2 \left| \int_{V_1} e^{i\mathbf{q}\cdot\mathbf{r}} d\mathbf{r_1} \right|^2$$

Principle of Babinet: Same coherent scattering of these samples



### Contrast variation and contrast matching





# Structure & Form Factor



$$\frac{d\Sigma}{d\Omega}(\mathbf{q}) = \frac{1}{V}(\rho_1 - \rho_2)^2 \left| \int_{V_1} e^{i\mathbf{q}\cdot\mathbf{r}} d\mathbf{r_1} \right|^2$$



 $\frac{d\Sigma}{d\Omega}(q) = \frac{N}{V}(\rho_1 - \rho_2)^2 V_p^2 P(q) S(q)$ 



Split up the integral over the sample



Fig. 2. Observed (•) and calculated (—) scattered intensity I(Q) as a function of momentum transfer Q for a charged micellar dispersion: 0.03 mol dm<sup>-3</sup> hexadecyltrimethylammonium chloride in D<sub>2</sub>O at 313 K. The functions P(Q) and S(Q) are discussed in the text. (1 barn sterad<sup>-1</sup> = 10<sup>-28</sup> M<sup>2</sup> sterad<sup>-1</sup>). Form factor P(Q) Interference of neutrons scattered at the same object

Shape, surface and densitiy distribution of objects

Structure factor S(Q)

Interference of neutrons scattered from different objects

> Arrangement or superstructure of objects

Measurement signal: Convolution of P(Q) and S(Q)







# a very diluted sample of randomly distrib

Lets assume a very diluted sample of randomly distributed identical scattering objects: No structure factor!

Form factor P(Q)

Interference of neutrons scattered at the same object

z

Patterson function: Convolution of an object with itself  $P(\mathbf{r}) = \int \rho(\mathbf{r}_1)\rho(\mathbf{r}_2) \ d\mathbf{V}, \quad \text{with } \mathbf{r} = \mathbf{r}_1 - \mathbf{r}_2$ 

Characteristic function (2D): Orientational average of P(r)  $\gamma(r) = \frac{1}{2\pi} \int_{0}^{2\pi} P(\mathbf{r}) d\varphi$ 

FT of Patterson function: Scattering signal

 $I(Q) = 4\pi \int_0^D \gamma(r) \frac{\sin(Qr)}{Q} r \,\mathrm{d}r$ 









Form Factor



## Form Factor





The form factor yields information about: Shape, size (distribution), SLD contrast (profile) and volume concentration of scattering objects.

#### But: Prone to overfitting due to <u>lack of independent</u> <u>parameters</u>!



Form Factor



Low and high Q limits: Porod and Guinier regime



Porod scattering for smooth surfaces and  $Q \gg 1/D$ 

 $I(q) \propto (q)^{-4}$  $\frac{\pi}{Q^{\star}} \cdot \lim_{q \to \infty} (I(q) \cdot q^4) = \frac{S}{V}$ 

Guinier scattering for dilute, monodisperse and isotropic solutions of particles:  $QR_G \ll 1$ 

$$I(q) = I(0)e^{\frac{-(qR_g)^2}{3}}$$





**FRM II** 

Forschungs-Neutronenquelle

Heinz Maier-Leibnitz

collimation length (1m; 1.5m; 2m; 3m; 4m; 6m; 8m; 10m; 12m; 14m; 16m; 18m; 20m)



FRM II Forschungs-Neutronenquelle Heinz Maier-Leibnitz

# **Collimation section**





Velocity selector



- Collimation: Define resolution and intensity
- Aperture system/neutron guides (supermirror)
- Alignment extremely critical
- Well-defined and homogenous wavelength /divergence profiles
- Transmission polarizer for the use of polarized neutrons
- Parasitic background scattering has to be avoided (whole collimation system is evacuated, edge scattering, incoherent background)



## Sample stage



Provide necessary sample environment (similar to any other neutron diffractometer or spectrometer)

Parasitic background scattering has to be avoided (extremely critical!)



Minimize neutrons travelling in air (few cm can be too much)

Avoid Aluminum neutron windows (single crystalline sapphire is better)

Get rid of scattering at edges (use conical slits)





## Detector tube



#### Vacuum vessel for detector to provide lowest possible background





Sample detector length adjustable (select Q-range)

One (or several) He<sup>3</sup> position sensitive detectors (typical 1m<sup>2</sup> with 5mm resolution) Typical length 10-40m

Interior completely covered with neutron absorbing Cadmium



Detector







What is the largest object SANS can detect (limit small Q)?

For large scattering angles (large Q) wavelength resolution dominates.

For 
$$a_1 = a_2 = a$$
 and  $L_1 = L_2 = L$   
 $\Longrightarrow \delta Q \approx \frac{\delta \theta}{\theta_{min}} Q_{min} \approx \delta \theta \frac{4\pi}{\lambda} \approx \frac{2\pi a}{\lambda L}$   
Largest object:  $\frac{2\pi}{\delta Q} = \frac{\lambda L}{a}$  On D11, ILL: L=40m,  $\lambda$ =15Å  $\Longrightarrow$  D $\approx$  5µm



Small Angle Neutron Scattering, S. Mühlbauer 19.10.2020







Intensity: Quadratic decrease with source to sample distance (collimation length)

Wavelength: Decrease of intensity with  $\lambda^{-4}$ 

$$\delta Q \approx \frac{\delta \theta}{\theta_{min}} Q_{min} \approx \delta \theta \frac{4\pi}{\lambda} \approx \frac{2\pi a}{\lambda L}$$

# **Typical Measurement**



Typical SANS dataset: •Sample (at different L) •Water (absolute scale, solid angle correction) •Empty sample holder/cuvette •Background (empty beam/blocked beam)

Forschungs-Neutronenquelle

Heinz Maier-Leibnitz

**FRM II** 

Fit model of the sample (conv. with resolution to the dataset)









### SANS: Diffractometer specialized for small scattering angles

Large correlations in real space 20 to 40000Å

Low **Q** small scattering angles  $\sim 1 \text{ Å}^{-1}$  to  $\sim 10^{-4} \text{ Å}^{-1}$ 

SANS tells you the statistical average of:

- Shape of scattering object
- Size (distribution) of scattering objects
- Surface of scattering objects
- Scattering length density (distribution)
- Arrangement (Superstructure?)











### Applications & Examples of SANS:

- Soft Matter
- Hard Matter
- Magnetism of nanoscale materials
- Emergent nanoscale magnetic structures of strongly correlated electron systems





Polyisoprene – Polystyrene Diblock Copolymer Phase Diagram

Two (or more) homopolymers units linked by covalent bonds Microphase separation: Complex nanostructures phases Flory-Huggins segment-segment interaction Degree of polymerization Volume fraction





**Figure 10.** Contour plots of SANS patterns for the sample IS-68 in three different ordered states (A, D, B). A nearly featureless pattern (a) was observed in state A (120 °C), which is attributed to a parallel orientation of lamellae with respect to the shear plane. After heating the sample to 145 °C (state D) and application of dynamic shearing ( $\dot{\gamma} = 0.1 \text{ s}^{-1} \text{ with } |\gamma| = 300\%$ ), a weak hexagonal scattering pattern was observed (b). This is consistent with a hexagonal in-plane arrangement of perforations in the minority (PS) layers. Further heating the sample to 210 °C, without shear (state B), produced a predominantly four-peak pattern (c), which transformed into result (d) when a shear rate of 2.2 s<sup>-1</sup> was applied. The azimuthal relationship between the combined 10 reflections in (c) and (d) is consistent with the SANS data reported for sample IS-39 (ref 22) and the *Ia*3*d* space group symmetry.

Macromolecules, Vol. 28, No. 26, (1995) Physics Today, p. 32, Feb. (1999)

Spheres

Cylinders

Bicontinuous

Perforated Layers

Lamellae





### Applications & Examples of SANS:

- Soft Matter
- Hard Matter
- Magnetism of nanoscale materials
- Emergent nanoscale magnetic structures of strongly correlated electron systems



## Hard Matter



Co-Re alloys for ultra high temperature / high stress applications (turbine blades) Precipitation growth and ageing by means of SANS.





FIGURE 2.1.2: Micrograph from Co-Re-Cr-Ta-C alloy showing various carbide morphologies. (a) The lamellar  $Cr_{23}C_6$  type carbides are present with their orientation relationship to the matrix. Inset: Area diffraction pattern. (b) Fine dispersion of TaC, embedded in  $\epsilon$ -Co matrix, between the  $Cr_{23}C_6$  lamellae. The inset shows a TEM image of the TaC precipitates. Adapted from [17]



Small Angle Neutron Scattering, S. Mühlbauer 19.10.2020





### Applications & Examples of SANS:

- Soft Matter
- Hard Matter
- Magnetic SANS
- Magnetism of nanoscale materials
- Emergent nanoscale magnetic structures of strongly correlated electron systems

FRM II Forschungs-Neutronenquelle Heinz Maier-Leibnitz



Magnetism on the nanoscale (30Å -3000Å)

Magnetism of materials that are inhomogneous on the nanoscale.

Nanoscale magnetism is imprinted by imhomogeneous material/sample properties.

e.g.: Nanoparticles, ferrofluids, sintered NeFeB magnets, granular magnetic recording media, nanodots, magnetic steels....etc



Emergent nanoscale magnetism of spatially homogeneous materials.

Nanoscale magnetism emerges due to (emergent) electronic correlations of SCES, competing interactions or superconductivity

e.g.: Superconducting VL, skyrmions, complex charge / stripe order due to competing FM/AF interactions....









Magnetic SANS is sensitive to the deviation of the mean magnetization:

Nuclear grain microstructure







Spatial variation of magnetization

Dissertation D. Mettus, Univ. Luxembourg

Deviation of the mean spin direction
 Deviation of the mean magnetitude of the magnetization
 SANS: 3D Fourier transform of the magnetization distribution

$$\mathbf{M}(\mathbf{r}) = \{M_x(\mathbf{r}), M_y(\mathbf{r}), M_z(\mathbf{r})\} \\ = \frac{1}{(2\pi)^{3/2}} \int_{-\infty}^{+\infty} \int_{-\infty}^{+\infty} \int_{-\infty}^{+\infty} \widetilde{\mathbf{M}}(\mathbf{q}) e^{i\mathbf{q}\mathbf{r}} d^3\mathbf{q}$$

Attention: "Nuclear" signal in present additionally!





#### Goal: Relating the microstructure and magnetic properties is of high interest.

Unpolarized magnetic SANS cross-sections

$$\frac{d\Sigma_{\perp}}{d\Omega} = K \left( b_H^{-2} |\widetilde{N}|^2 + |\widetilde{M}_x|^2 + |\widetilde{M}_y|^2 \cos^2 \theta + |\widetilde{M}_z|^2 \sin^2 \theta - CT_{yz} \sin \theta \cos \theta \right).$$
$$\frac{d\Sigma_{\parallel}}{d\Omega} = K \left( b_H^{-2} |\widetilde{N}|^2 + |\widetilde{M}_x|^2 \sin^2 \theta + |\widetilde{M}_y|^2 \cos^2 \theta + |\widetilde{M}_z|^2 - CT_{xy} \sin \theta \cos \theta \right).$$

Use magnetic selection rule (anisotropic scattering pattern)

Saturate the sample with high fields to surpress magnetic scattering

Use SANSpol (half polarized) and POLARIS (full polarization analysis) for a full picture/ discrimination of SF/NSF scattering (comes at the price of intensity).



Fig. 2.12 (upper row) (*a*)-(*d*) Contour plots of normalized  $d\Sigma_M/d\Omega$  (Eq. (2.58)) at applied magnetic fields as indicated ( $\mathbf{k}_0 \perp \mathbf{H}_0$ ;  $H_p/\Delta M = 1$ ;  $\mathbf{H}_0$  is horizontal). For  $h^2(qR)$  and  $\widetilde{M}_z^2(qR)$ , we used the form factor of the sphere with a radius of R = 5 nm (Eq. (2.78); S(q) = 1). (lower row) (*e*)-(*h*) Corresponding two-dimensional correlation functions c(y,z), which were computed according to Eq. (2.91) ( $H_p/\Delta M = 1$ ).





SANS measures inhomogeneities of scattering length density

$$\frac{d\Sigma}{d\Omega}(\mathbf{q}) = \frac{1}{V}(\rho_1 - \rho_2)^2 \left| \int_{V_1} e^{i\mathbf{q}\cdot\mathbf{r}} d\mathbf{r_1} \right|^2$$

SANS measures mesoscopic information, independent of microscopic structure



FRM II Forschungs-Neutronenquelle Heinz Maier-Leibnitz

## **Incommensurate Magnets**



Archetypal helimagnet MnSi

Itinerant ferromagnet

No inversion symmetry (B20) Dzyaloshinsky-Moriya interaction

 $H=J(S_1 \cdot S_2) + D \cdot (S_1 \times S_2)$ 

Monochiral helix, pitch=180 Å.

Cubic anisotropy: Weak pinning (111)



# SANS: Incommensurate satellites around (0,0,0)



Diffraction: Incommensurate satellites around (h,k,l)







Vortex lattice crystallography: Diffraction in SANS geometry

Vortex lattice 2D magnetic Bravais lattice

$$\phi_0 = \frac{h}{2e} |\vec{a}_i| = \left(\frac{2\phi_0}{\sqrt{3}B}\right)^{\frac{1}{2}} |\vec{a}_i| = \frac{2\pi}{|\vec{Q}|}$$



Typical values:Intensity Bragg peakForm factorB=1500 G $R=\frac{2\pi \gamma^2 \lambda_n^2 t}{16\phi_0^2 Q}|h(Q)|^2$  $h(Q)=\frac{\phi_0}{(2\pi\lambda)^2}e^{\frac{-\pi B}{B_{c2}}}$  $A_0=1260 \text{ Å}$  $R=\frac{2\pi \gamma^2 \lambda_n^2 t}{16\phi_0^2 Q}|h(Q)|^2$  $h(Q)=\frac{\phi_0}{(2\pi\lambda)^2}e^{\frac{-\pi B}{B_{c2}}}$ 







## Thank you for your attention!



## www.mlz-garching.de Sebastian.Muehlbauer@frm2.tum.de